

## Solutions Atomic Physics 2

1. Probability distributions without any zero points ( $r > 0$ ), i.e. wavefunctions with  $n - \ell - 1 = 0$  zero crossings, correspond to orbitals with the maximum  $\ell$  for given  $n$ , i.e. 1s, 2p, 3d, 4f, 5g,.... For these orbitals it so happens that the radial position of the maximum corresponds to the radius of a circular Bohr orbit with the same main quantum number  $n$ :

$$r = a_0 \cdot \frac{n^2}{Z}.$$

From the figure we find the maximum at  $r = 16a_0$  and thus  $n = 4$ .  
Hence the orbital is 4f.

2. The center-of-gravity energy for  $4p^2P$  is

$$E = \frac{1}{6}(2 \cdot 12985.2 + 4 \cdot 13042.9) = 13023.7 \text{ cm}^{-1}$$

The term value is then  $T = E_{\text{limit}} - E = 21986.1 \text{ cm}^{-1}$  and the quantum defect

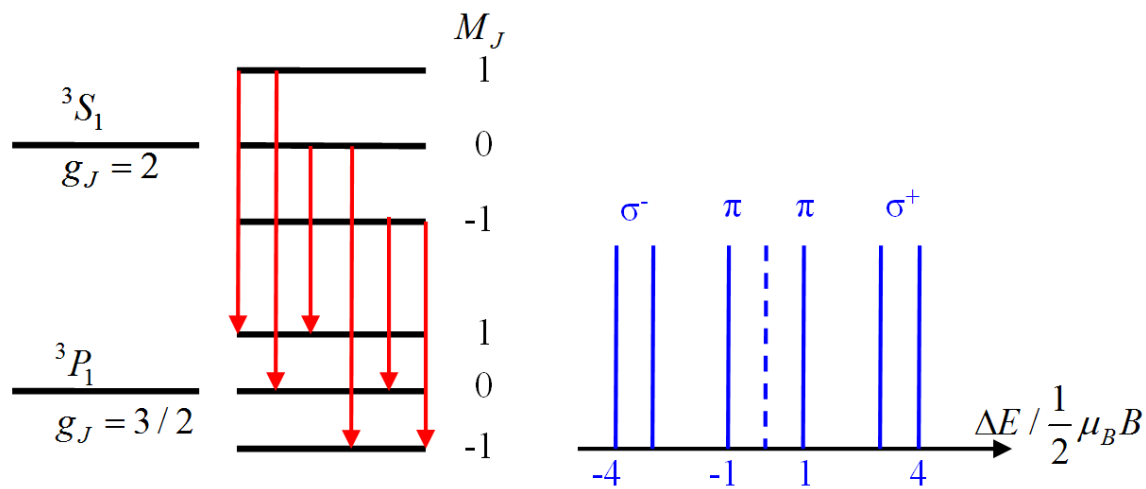
$$\delta = 4 - \sqrt{\frac{R_\infty \cdot 1^2}{T}} = 1.766.$$

Assuming  $\delta$  to be constant in the  $p$ -series we get

$$T(20p^2P) = \frac{R_\infty \cdot 1^2}{(20 - \delta)^2} = 330.0 \text{ cm}^{-1} \Rightarrow E(20p^2P) = 34679.7 \text{ cm}^{-1} \Rightarrow \lambda(4s - 20p) = 2883.5 \text{ \AA}.$$

The experimental energy of  $20p$  is  $34681.7 \text{ cm}^{-1}$ .

- 3ab. The Zeeman components in the  $^3P_1 - ^3S_1$  transition are shown below with their polarization state, when viewed viewed perpendicularly, indicated. Note that  $0 \rightarrow 0$  is not possible since  $\Delta J = 0$ .



3c. The minimum energy separation is

$$\Delta E_{\min} = \frac{1}{2} \mu_B B \Rightarrow \Delta f_{\min} = \frac{\mu_B B}{2h} = 3.5 \text{ GHz} \Rightarrow \Delta \lambda_{\min} = \frac{\lambda^2}{c} \Delta f_{\min} = 0,027 \text{ \AA}.$$

4a. If no other effects limit the experimental resolution the limit is set by the natural line

$$\text{width. } \Delta f_N = \frac{1}{2\pi} \left( \frac{1}{\tau_1} + \frac{1}{\tau_2} \right) = \frac{1}{\pi\tau} \text{ if the 2 lifetimes are equal.}$$

With  $\Delta f_{\min} = 3.5 \text{ GHz}$  we get  $\tau > 91 \text{ ps}$ .

(Using  $\Delta \lambda_{\min} = 0,05 \text{ \AA} \Rightarrow \Delta f_{\min} = 6.5 \text{ GHz} \Rightarrow \tau > 48.9 \text{ ps}$ .)

4b. If the limit is set by the Doppler width instead we get:

$$\Delta f_D = C \cdot \frac{c}{\lambda} \cdot \sqrt{\frac{T}{M}} \Rightarrow T < M \cdot \left( \frac{\Delta f_{\min} \cdot \lambda}{c \cdot C} \right)^2 =$$

$$112 \text{ g/mol} \cdot \left( \frac{3.5 \cdot 10^9 \text{ s}^{-1} \cdot 4800 \cdot 10^{-10} \text{ m}}{3 \cdot 10^8 \text{ m/s} \cdot 7.16 \cdot 10^{-7} (\text{g}/(\text{mol} \cdot \text{K}))^{0.5}} \right)^2 = 6851 \text{ K}$$

$$\Delta \lambda_{\min} = 0.05 \text{ \AA} \Rightarrow T = 23600 \text{ K instead.}$$

5a. From the radiation balance:  $(A_{21} + \rho(f_{21})B_{21}) \cdot N_2 = \rho(f_{21})B_{12} \cdot N_1$  we find that if the number of stimulated photons should exceed the number of absorbed we must have:

$$\rho(f_{21})B_{21} \cdot N_2 > \rho(f_{21})B_{12} \cdot N_1 \Leftrightarrow \frac{g_1}{g_2} B_{12} \cdot N_2 > B_{12} \cdot N_1 \Leftrightarrow g_1 N_2 > g_2 N_1.$$

If  $g_1 = g_2$  this reduces to the more direct condition of an inverted population  $N_2 > N_1$ .

5b. When the atom absorbs a photon that carries a momentum  $p = h/\lambda$  it must change its velocity ( $\Delta v$ ) so that the total momentum is conserved.

$$\bar{p}_\gamma + \Delta \bar{p}_{\text{atom}} = 0, \text{ i.e. the velocity change by: } \Delta v \cdot M = \frac{h}{\lambda}.$$

6.  $\ell_1 = 1, s_1 = 1/2 \Rightarrow j_1 = 1/2, 3/2$ .

$$j_1 = 1/2, \ell_2 = 2 \Rightarrow K = 3/2, 5/2. \quad j_1 = 3/2, \ell_2 = 2 \Rightarrow K = 1/2, 3/2, 5/2, 7/2.$$

If we now add  $s_2 = 1/2$  we will get levels with  $J = K \pm 1/2$ .

$$\frac{1}{2}[3/2]_{1,2}, \frac{1}{2}[5/2]_{2,3}$$

$$\frac{3}{2}[1/2]_{0,1}, \frac{3}{2}[3/2]_{1,2}, \frac{3}{2}[5/2]_{2,3}, \frac{3}{2}[7/2]_{3,4}.$$

(Note that we get the same *number* of levels and the same set of total  $J$ : s as in  $LS$ - or  $jj$ -coupling)

This coupling scheme could be a good approximation in a case where the largest interaction is the spin-orbit of the most tightly bound (non-*s*) electron (2*p*), followed by the electrostatic repulsion between the 2 electrons and finally the very small spin-orbit contribution of the outer, high-*l* electron (5*d*). Remember that the spin-orbit energy scales as  $n^{-3}$  and  $l^{-3}$ .

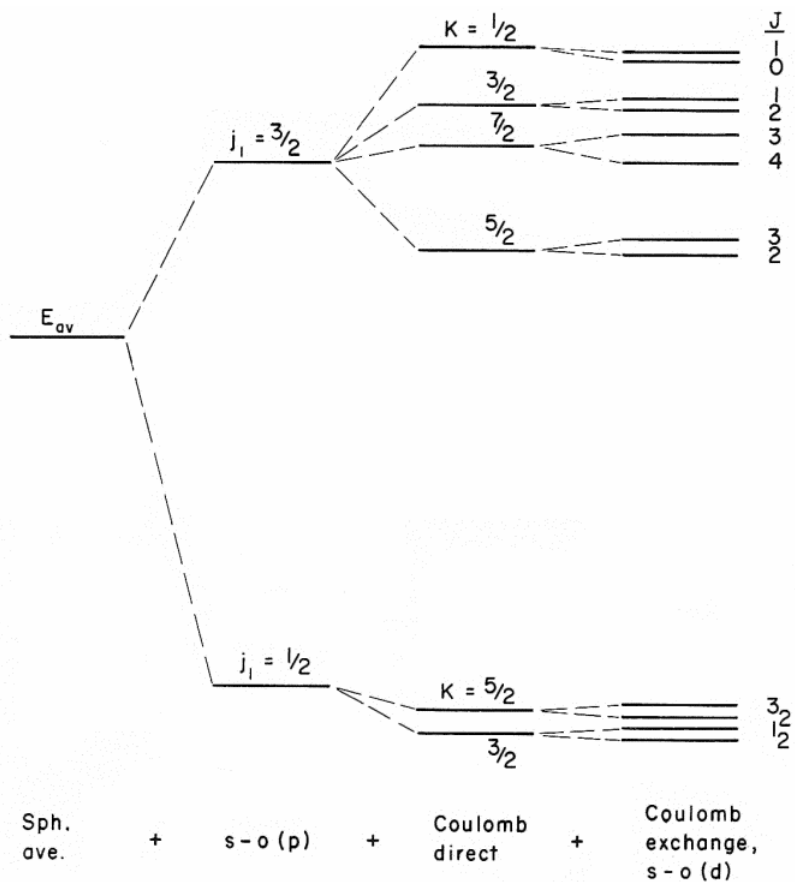


Figure: *pd*-configuration in *jK*-coupling